### Research Article

# Wenzheng Lu\*, Leonardo de S. Menezes, Andreas Tittl, Haoran Ren and Stefan A. Maier\* Active Huygens' metasurface based on *in-situ* grown conductive polymer

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Abstract: Active metasurfaces provide unique advantages for on-demand light manipulation at a subwavelength scale for emerging visual applications of displays, holographic projectors, optical sensors, light detection and ranging (LiDAR). These applications put stringent requirements on switching speed, cycling duration, electro-optical controllability, modulation contrast, optical efficiency and operation voltages. However, previous demonstrations focus only on particular subsets of these key performance requirements for device implementation, while the other performance metrics have remained too low for any practical use. Here, we demonstrate an active Huygens' metasurface based on conductive polyaniline (PANI), which can be insitu grown and optimized on the metasurface. We have achieved simultaneously on the active metasurface switching speed of 60 frame per second (fps), switching duration of more than 2000 switching cycles without noticeable

degradation, hysteresis-free controllability over intermediate states, modulation contrast of over 1400 %, optical efficiency of 28 % and operation voltage range within 1 V. Such PANI-powered active metasurface design can be readily incorporated into other metasurface concepts to deliver high-reliability electrical control over its optical response, paving the way for compact and robust electrooptic metadevices.

Keywords: active metasurfaces; electrical switching; conductive polymer; nanoantenna; beam steering

### **1** Introduction

Active metasurfaces, often alternatively termed as tunable or reconfigurable metasurfaces, are rapidly emerging as a major frontier in photonic research and have launched tremendous breakthroughs in modern optics [1]-[3]. Compared with their passive counterparts, active metasurfaces consist of ultrathin planar arrays of subwavelength active nanoantennas, whose optical responses can be dynamically modulated on-demand. Over the past decade, active metasurfaces have made a significant impact on the new development of beam deflectors [4]-[8], spatial light modulators [9]-[12], varifocal metalenses [13]-[15], dynamic holograms [16]-[19] and many others. The active tuning schemes mainly rely on varying optical properties of the nanoantennas or their surrounding materials through chemical reactions [19]-[21], mechanical displacements [22]-[24], electrical switching [16], [17], [25]-[27], thermal modulation [28]-[31] and all-optical switching [32]-[34]. Among these modulation schemes, electrical switching is of particular interests because it promises compact integration of meta-optics with miniaturized on-chip electro-optic systems, which can be readily incorporated into electronic smart devices for practical applications [1], [2], such as displays, augmented/virtual reality (AR/VR) glasses, dynamic holograms, optical sensing, beam steering and light detection and ranging (LiDAR).

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To this end, a number of active metasurfaces have been implemented by using different electro-active materials, such as chalcogenide phase-change materials (PCMs) (26, 27), III-V semiconducting materials [5], [35], [36], ionic conducting materials [37], [38], metallic polymers [6], [14]. Previous active metasurfaces exhibit only distinct subsets of key performance metrics [1], including switching speed, cycling duration, electro-optical controllability, modulation contrast, optical efficiency and operation voltages, as shown in Figure 1a. However, these performance attributes are hardly met simultaneously, making them almost impossible for any practical use. For instance, even though PCM metasurfaces can provide a switching speed of up to 2 MHz [27], they require high operation voltage and large cell manipulation [26]. III–V semiconductors offer excellent switching duration and controllability for intermediate states, but suffer from low modulation contrasts due to the small tuning range of the material intrinsic refractive index [5]. Recently, conductive polymers have shown many desired properties for electrically active metasurfaces, such as large variation of refractive index, fast switching speed, superior cycling stability and low operation voltages [39]-[42]. However, previous active metasurfaces purely based on conductive polymers are disadvantaged by low diffractive efficiency, due to the geometric phase design and the weak resonant nature of the nanostructured polymer [6], [14]. Moreover, the fabrication of the conductive polymer relying on spincoating prevents long-term switching durability since the contact issue between the polymer and the substrate. While active metasurfaces incorporating conductive polymers as active surrounding materials have shown superior switching performance [16], the implemented geometric phasebased plasmonic nanoantenna design yields suboptimal diffraction efficiency in optical frequencies.

In this work, we introduce an electrically active Huygens' metasurface based on the in-situ grown conductive polymer, polyaniline (PANI), and experimentally demonstrated the overall active performance. We combine the superior electro-optical response of PANI with the dielectric Huygens' nanoantennas by an electrochemically in-situ grown manufacture methods, allowing for the strengthened cycling stability and the optimization of electro-optical modulation. Our active Huvgens' metasurface exhibits a high modulation contrast of 1400 % and a diffraction efficiency up to 28 %, which is 25 times higher than the previous polymer-based metasurface [6]. The solid contact between the nanoantennas and the grown PANI facilitates mechanical and electrical durability, enabling a superior cycling duration of over 2000 switching cycles without noticeable degradations. The intrinsic dynamic

properties of PANI endow the active metasurface with fast switching speed of 60 fps and hysteresis-free controllability within a low operation voltage range from -0.2 V to +0.6 V. Unlike previous active metasurfaces performing well only in subsets of performance metrics (Figure 1a), our active metasurface holds great promise towards practical applications that needs overall switching performance.

# 2 Concept of electrically active Huygens' nanoantennas

Our active Huygens' metasurface is composed of dielectric silicon nanodisks surrounded by a layer of PANI, resting on an indium-tin-oxide (ITO) substrate that acts as an electrical contact (Figure 1b). The cylindrical shape of the nanoantenna is chosen to provide polarization independence and the period of the nanoantennas is fixed to 450 nm in both x and y direction so that the unit cell is subdiffractive. The state of the polymer is collectively controlled by voltage tuning via electrochemistry, where an applied voltage range between +0.6 V and -0.2 V (vs. a reference electrode) can continuously tune the polymer into states in between the oxidized and reduced states. In detail, the redox reaction triggered by the applied voltage induces an alternating delocalization of the  $\pi$ -electrons in the polymer chain of PANI, resulting in a large variation of the refractive index with a maximum  $\Delta n = 0.6$  at 780 nm (Figure S1, Supporting Information).

The core concept of our design is the spectral tuning of the electric dipole (ED) and magnetic dipole (MD) of the nanoantenna by electrical modulation of the surrounding polymer refractive index. Figure 1c and d depict the spectral response of nanoantennas at different applied voltages. In the PANI oxidized state, the simulated and measured metasurface transmission spectra show that the ED and the MD are spectrally separated, resonating at 694 nm and 765 nm, respectively. When the polymer is switched to the reduced state, the two dipole resonances redshift to 785 nm and overlap with each other. As suggested by the principle of dielectric Huygens' metasurfaces [43], the spectrally overlapping resonances of ED and MD signifies the fulfillment of Huygens' condition, which produces a full-range phase shift from 0 to  $2\pi$  combined with high transmission. Importantly, the phase modulation functionality of the nanoantenna is switched on at an applied voltage of -0.2 V, exhibiting strong phase shift (Figure 1d), while at +0.6 V, the nanoantenna exhibits a phase retardation well below  $2\pi$ -phase coverage and can be considered switched off (Figure 1d). The variation of the phase shift in the nanoantennas is



**Figure 1:** Comparison of active metasurface performance and the concept of electrically active Huygens' metasurface composed of PANI-modulated dielectric nanoantennas. (a) Comparison of different electrically active metasurfaces in performances. The performance is evaluated relatively as defined in the reference [1]. The evaluation of each electro-active materials is based on representative references: chalcogenide PCMs [26], [27], III–V semiconductors [5], [36], ionic conducting materials [37], [38], metallic polymers [6]–[14]. (b) Schematics illustrating the electrical switching of an individual active Huygens' nanoantenna based on conductive polymer. The polymer is switched to oxidized state at +0.6 V (left), and to reduced state at -0.2 V (right). The switching time between the states is 15 ms, corresponding to a refresh rate of over 60 fps. (c) Simulated and measured transmission spectra (left panel) of a metasurface made of homogeneous nanoantennas (height H = 140 nm, radius R = 110 nm, periodicity  $P_x = P_y = 450$  nm) and the calculated phase retardation (right panel) of the nanoantennas with different radii in the oxidized state. (d) Simulated and measured transmission spectra (left panel) and the calculated phase retardation (right panel) in the reduced state. The inset chemical structures show the conversion of the polymer between the oxidized state, as well as the spectrally overlapping electric and magnetic dipoles (ED + MD, yellow) in the reduced state, at their corresponding resonant wavelengths of 694 nm, 765 nm and 785 nm, respectively.

determined by the interference of the ED and MD resonances, where the spectral responses of the two resonances are determined by the nanoantenna geometry and the refractive-index contrast between the nanoantenna and the surrounding environment [43]. Therefore, by harnessing the environmental tuning scheme and carefully choosing nanoantenna geometries for optimum phase engineering, switchable metasurfaces with high efficiency and greatly improved switching contrast can be realized.

# 3 Metasurface design for electrical beam steering

As a proof of concept, we implemented an active metasurface for optical beam steering, where the incident light can be deflected into a fixed angle on-demand by electrical control. The metasurface design is based on an eight-element gradient phase geometry using the calculated nanoantenna phase profiles at the reduced state for an incident wavelength of 785 nm (Figure S2). This design constructs a modulator between binary states. It allows for transmitted beam deflection at the reduced state, while maintains the undeflected beam at the oxidized state (Figure 2a and b). The chosen structures have an average transmittance of 0.6. The metasurfaces were fabricated via electronbeam lithography (EBL) patterning and reactive-ion etching (RIE), followed by electrochemical polymerization for the growth of the polymer. Figure 2c shows a scanning electron microscopy (SEM) image of the fabricated bare dielectric metasurface on a transparent conductive ITO/glass substrate. The metasurface is designed such that the steered beam couples to the +1 diffraction order at the polymer reduced state, and to the 0 order at the polymer oxidized state. The design principle is validated by simulated electric field distributions (Figure S3), where the transmitted spatial phase profile supports diffraction at +1 order at an operation voltage of -0.2 V and strongly suppresses the diffraction at +0.6 V. Detailed far-field analysis of the metasurfaces reveals the switchable beam steering performance with high switching contrast at a diffraction angle of 12.7°, meanwhile the intensity of the 0 order can be readily modulated (Figure 2d).

# 4 *In-situ* polymer growth and electrical beam steering performance

To experimentally validate our electrical beam steering metasurfaces, electrochemical growth of PANI was implemented on the fabricated metasurface *in-situ* in the optical measurement setup, where a high-speed



**Figure 2:** Metasurface design for electrical beam steering. (a – b) Schematic illustrating the concept of electrical beam steering using a polymer-integrated Huygens' metasurface at the applied voltage of +0.6 V and -0.2 V, respectively. (c) Tilted SEM image of the bare dielectric metasurface with eight-element gradient nanoantennas design (highlighted structures). Scale bar: 1  $\mu$ m. (d) Simulated far-field transmitted intensities of +1, 0, -1 diffraction orders for the electrical beam steering metasurface.

monochromatic camera was used to monitor the diffraction pattern during the polymer growth (Figure 3a). A threeelectrode system integrated with the optical measurement setup was employed for PANI growth and electrical switching in an aqueous electrolyte using a custom-built electrochemical cell (Figure S4, see also Section 7 for further information). Since PANI is switched between -0.2 V and +0.6 V and can be electrochemically polymerized under an applied voltage of +0.8 V on an ITO electrode [44], [45], a cyclic voltammetry with a voltage range from -0.2 V to +0.8 V was chosen to carry out PANI growth and electrical switching simultaneously (Figure S5, see also



**Figure 3:** *In-situ* measurement and optimization of beam steering performance. (a) Schematic of the *in-situ* optical measurement setup for the active beam steering metasurface. A high-speed monochromatic camera was used to record the diffraction pattern during the polymer growth process. (b) Normalized intensities of +1, 0 and -1 diffraction orders during the polymer growth. (c) Figure of merit (FOM) of beam steering between 0 and +1 diffraction orders at different coating cycles obtained from the measured intensities in (b). (d) Measured camera image (top) and diffraction intensity profile (bottom) at an applied voltage of +0.6 V. (e) Measured camera image (top) and diffraction intensity profile (bottom) at an applied voltage of -0.2 V. (f) Camera images of the diffraction orders at different input voltages. (g) Diffraction intensity profiles at different input voltage extracted from (f). The data in (d)-(f) were measured on the beam steering metasurface at an optimized polymer coating cycle number of 56.

Supplementary Note 1 in Supporting Information), thus allowing for *in-situ* optimization of the electrical beam steering performance.

The diffraction images were monitored throughout the entire polymer growth process. Figure 3b shows the intensities of +1, 0 and -1 diffraction orders as function of the PANI coating cycle. In the first 25 coating cycles, the variation of the diffractive intensities is very small due to the insufficient amount of the grown PANI. By increasing the coating cycle number and thus the PANI thickness, the effect of the grown PANI layer becomes large enough to enable switching of the diffraction intensities. As predicted by the design, the diffraction intensity of the +1 order reaches its minima and maxima at the oxidized state and the reduced state, respectively, while the 0 order shows an opposite response. The intensity contrasts of the +1 and 0 diffraction orders at the two states are dictated by the optical phase variation and the intrinsic optical absorption of PANI. Further increasing the PANI thickness allows for the optical absorption of PANI to become more influential, as evidenced by the decreased diffractive intensity and switching contrast in the +1 order. The overcoated PANI also leads to a phase mismatch in the nanoantennas, where the switching of the +1 and 0 diffraction orders becomes synchronous and solely dependent on variation of optical absorption of PANI.

Quantitative evaluation of the electrical beam steering performance is accomplished by a figure-of-merit (FOM) generally used for beam deflectors, which is based on the switching contrasts between the +1 and 0 diffraction orders at two states [28]. The FOM can be represented as:

$$FOM = \frac{1}{2} (C^{ox} - C^{red})$$
$$= \frac{1}{2} \left( \frac{I_0^{ox} - I_1^{ox}}{I_0^{ox} + I_1^{ox}} - \frac{I_0^{red} - I_1^{red}}{I_0^{red} + I_1^{red}} \right)$$

where *C* is the switching contrast between two diffraction orders,  $I_0$  and  $I_1$  represent the intensities of the 0 and +1 diffraction orders, respectively, and the superscripts of ox and red denote the oxidized and reduced states, respectively. The value of FOM, which falls between 0 and 1, describes the capability of a beam deflector to route energy into certain angles, where FOM = 1 represents the ideal case. The calculated FOM for different coating cycles indicates that the beam steering performance is optimized when the FOM reaches 0.6 with a polymer coating cycle number of 56 (Figure 3c). In comparison, the obtained FOM from a chalcogenide PCM-based beam steering metasurface [26] is 0.42, and 0.38 from a liquid crystal-based one [9]. Figure 3d and e present the diffraction images at the optimized thickness and the corresponding intensity profiles, at +0.6 V and -0.2 V, respectively, and the results are in excellent agreement with the simulation. The modulation contrast defined as the intensity ratio of the +1 order between the two states, is over 1400 % and can be further increased by increasing the PANI thickness. The thickness of the PANI laver at the optimized condition is measured to be 200 nm (Figure S6), which is only 1/7 as thin as the thickness of the embedding active layer in liquid crystal-based active metasurfaces [9], [17]. The PANI layer grown on the metasurfaces has completely covered the entire resonators with a fiber-like surface structure, as shown in the SEM image (Figure S7). So far, the diffraction efficiency of the +1 order at the reduced state is 28.8 %. In comparison with the beam steering metasurfaces purely based on conductive polymers [6], [14], our design has a 25 times higher diffraction efficiency. The optical loss increases when the PANI is switched to oxidized state, which accounts for up to 70 % of the transmitted energy of the reduced state (Figure S8). The diffraction efficiency can be further improved by decreasing the polymer thickness together with optimizations of the nanoantenna geometry and materials.

In addition, intermediate PANI states, which possess gradually changing refractive indices due to the partially delocalized  $\pi$ -electrons, can be accessed via voltage tuning. By operating the voltage between the oxidized and reduced states, we can continuously modify the diffraction intensity and the intensity ratio between the +1 and 0 diffraction orders, as demonstrated in the diffraction images (Figure 3f) and the intensity profiles (Figure 3g) at different applied voltages. Notably, the electrical switching demonstrates high-quality reversibility. With a cycling voltage between -0.2 V and +0.6 V, the electrical switching is fully reversible with a remarkable stability on our active metasurface, as shown in the Movie S1 which includes the diffraction images, diffraction intensities and the cyclic voltammograms for a total of 9 switching cycles in real-time. We further analyzed the hysteresis behavior of the electrical switching (Figure S9). Our electrically active metasurfaces can be operated in a nearly hysteresis-free manner. This allows switching to the intermediate states without a memory effect that could possibly deteriorate the accuracy of optical response by electrical modulation, which is crucial for realizing active metasurfaces for precise on-demand control of phase and amplitude.



**Figure 4:** Electrical switching speed and durability. (a – b) Temporal responses for the switch-on and switch-off processes, respectively. The rise time  $(\tau_{rise})$  of the switch-on process is 14.1 ms and the fall time ( $\tau_{fall}$ ) of the switch-off process is 11.7 ms. The rise time and the fall time are defined as the time required to switch between 10 % and 90 % of the optical response, depicted as the gray area, under an abrupt alternation of the applied voltage (green and red areas). (c) Input voltage schemes with different voltage duration times ( $\Delta t$ ) in a range from 50 ms to 5 ms. (d) Temporal optical response at different input voltage frequencies. Each scheme lasts 10 s. The percentage numbers indicate the maximum accessible range of the optical response. The recorded slightly inconsistent maxima and minima are due to frame drops of the camera. (e) Switching durability of the electrically active metasurface under  $\Delta t = 3$  s. The electrical switching is highly reversible and consistent without noticeable degradation over 2000 switching cycles. The optical response in this figure corresponds to the +1 order at an optimized coating cycle of PANI.

# 5 Electrical switching speed and durability

To evaluate the switching speed of our active metasurfaces, the temporal optical response is measured under an abrupt alternation of the input voltages between the on (-0.2 V)and off (+0.6 V) states, as presented in Figure 4a and b. The rise time ( $\tau_{rise}$ ) of the switch-on process and the fall time  $(\tau_{fall})$  of the switch-off process, defined as the time required for the modulated intensity to rise or fall between the 10 % and 90 % switching window, are 14.1 ms and 11.7 ms, respectively. The limiting factor of the switching speed is the electrical characteristics of the cell and the electrochemical system, where the potential switching time for polyaniline can be within 200 µs [47]. We further analyzed the switching stability under the high-speed and low-speed operation configuration. Under the high-speed working condition, the metasurface is operated at alternating input voltages between the on and off states with a voltage duration time ( $\Delta t$ ) varied from 50 ms to 5 ms (Figure 4c). Such intense switching conditions could possibly lead to mechanical or electrical failures due to the volume expansion, physical exfoliation or irreversible reactions of polymer films [46]. However, remarkably, when  $\Delta t$  approaches the switching time  $au_{\mathrm{rise}}$  and  $au_{\mathrm{fall}}$ , the accessible range of the optical response approaches the 10 %–90 % switching window at  $\Delta t = 15$  ms, equivalent to 66 fps, as shown in Figure 4d. Lowering  $\Delta t$ below 15 ms decreases the accessible optical range, but does not significantly influence the reversibility. In the low-speed switching condition, the metasurface is operated at  $\Delta t =$ 3 s and shows no decrease in signals within the operating window. The electrical switching also exhibits no substantial degradations in the switching contrast over 2000 switching cycles (Figure 4e). The key characteristic supporting the superior cycling stability is that, when electron injection or extraction takes place by voltage application, the highly conductive PANI is able to redistribute delocalized  $\pi$ -electrons along the polymer chain and keep mobility without structural degradation [47], [48], even after 10<sup>7</sup> switching cycles [49]. This ensures the stability when deploying intense highspeed electrical modulation, and the long-term switching contrast. Moreover, the solid contact between the nanoantennas and the PANI film, which is directly grown on the metasurface, provides additional mechanical and electrical stabilities.

### 6 Conclusions

In this study, we have introduced and experimentally realized an active Huygens' metasurface based on the insitu grown PANI with an overall switching performance. We have achieved electrical switching on the beam steering metasurface with large modulation contrast of over 1400 %, hysteresis-free electro-optical controllability, high diffraction efficiency of 28 %, fast switching speed of over 60 fps and superior switching duration over 2000 switching cycles, under an operation voltage range only from -0.2 V to +0.6 V. We have demonstrated that the active Huygens' nanoantenna design boosts the optical efficiency compared to the nanoantenna purely based on polymers. In addition, we have shown that PANI is an excellent candidate of electro-active materials for active metasurfaces, owing to its large tuning range of refractive index, fast switching speed, electrical properties and superior physical and chemical stability. The manufacture of conductive polymer by in-situ growth on metasurface is straightforward and more reliable, which can be processed in a low-cost, large-area and scalable manner. The concept of our electrically active Huygens' nanoantenna design can be broadly extended to a variety of active metasurfaces that are potentially capable of programmable and dynamic manipulation. From an industrial standpoint, the overall switching performance, together with the scalable, reliable, cost-effective polymer manufacturability, provide promising opportunities to achieve electro-optical metadevices for emerging applications, forging the path of active metasurfaces towards commercial success.

### 7 Materials and Methods

### 7.1 Numerical simulation

For the numerical simulation of the active Huygens' nanoantenna, we employed a finite-difference time-domain solver (Ansys Lumerical) using periodic boundary conditions in the x and y directions. The silicon material data took from Palik [50] and the polymer index took from the experimentally measured results obtained through ellipsometry (Figure S1). For the geometric parameters, the Si nanodisks have a height of 140 nm and a varying radius from 60 nm to 160 nm, and the height of polymer is 200 nm. Transmission spectra, phase and transmittance of the nanoantenna were simulated under a plane-wave excitation with a linear polarization and wavelength range from 500 nm to 1000 nm.

### 7.2 Metasurface fabrication

The dielectric metasurfaces were fabricated by a nanofabrication procedure that contains EBL patterning, mask deposition and RIE. In detail, a 140-nm thick Si film was deposited on a transparent conductive 50-nm thick ITO-coated fused silica glass substrate using a plasmonenhanced chemical vapor deposition. For the lithography, a double layered poly(methyl-methacrylate) (PMMA) positive photoresist (495 k A4 and 950 k A2) was spin-coated onto the Si film with a soft-baking for 90 s at 170 °C, followed by spin-coating of a conducting layer (ESpacer 300Z) on the photoresist to avoid electron charge and pattern distortion. The array of Si nanodisks was patterned on the photoresist using EBL (Raith eLine Plus), followed by a development process by immersing the sample into a 3:1 Isopropanol:Methylisobutylketone solution for 50 s. A 30-nm thick chromium layer was deposited using an e-beam evaporation as a hard mask. Lift-off process was carried out in a remover solution (Microposit Remover 1165). The designed pattern was finally etched into the Si film by a RIE process (Oxford Instruments) for 2 min. As the last step, the chromium hard mask was chemically removed by wet etching using a chromium-selective etchant solution (Sigma-Aldrich).

#### 7.3 Electrochemical setup

The electrochemical polymer growth and electrical switching were carried out in a specially custom-built electrochemical cell. The electrochemical cell is designed for housing a three-electrode system into a thin layer of aqueous electrolyte with an optical thickness of 1 mm on the top of the sample substrate, where a thin transparent glass was used to seal the electrolyte on the top, allowing for an optical transmission measurement through the metasurface. The sample substrate (ITO-coated fused silica) was used as the bottom sealing glass of the cell. The cell also features on the side the in-let and out-let for electrolyte to enable flow-in and flow-out for electrolyte replacement, which is driven by an electrical injector. The sample substrate was connected as the working electrode through a striped metal plate as the electrode contact, whereas a Pt wire and a Ag/AgCl reference were inserted from the side of the cell and connected as the counter electrode and the reference electrode, respectively. A potentiostat (CHI-760e) is used to apply voltage over the time to perform electrochemical polymer growth and switching.

#### 7.4 In-situ polymer growth and electrical switching

In-situ PANI growth was carried out by an electrochemical coating method, as reported previously [44]. The ITO substrate with the metasurface sample was connected as the working electrode. A cycling voltage in the range from -0.2 V to +0.8 V at a scanning speed of 25 mV/s was applied on the sample in an acidic aqueous electrolyte containing 1.0 M H<sub>2</sub>SO<sub>4</sub> and 0.2 M aniline. The thickness of the grown PANI thickness can be controlled by the number of voltage cycle.

For the electrical switching, the electrolyte in the electrochemical cell was replaced by an aniline-free aqueous electrolyte containing only 1.0 M  $H_2SO_4$ . To switch the PANI to the oxidized state and reduced state, constant voltages of +0.6 V and -0.2 V, respectively, were applied on the substrate. For electrical switching cycling, a cycling voltage in the range between -0.2 V and +0.6 V at a scanning speed of 25 mV/s was used.

#### 7.5 Structural characterization

Scanning electron microscopy (SEM) of the metasurface was performed on the Raith eLine Plus system in a SEM mode under a working voltage of 5 kV. The height of the deposited Si film and the thickness of the grown PANI layer at the optimized coating cycle were measured with a profilometer (Bruker Dektak XT) using a stylus with a radius of 2  $\mu$ m.

#### 7.6 Optical characterization

The refractive index of PANI was obtained from an ellipsometry measurement on an ellipsometer with dual-rotating compensators and a spectrometer (J.A. Woollam, M2000XI-210). A 100-nm thick PANI film electrochemically prepared on an ITO-coated glass substrate was measured with angle-variable spectroscopic ellipsometry at incident angles of 65°, 70° and 75°, using a bare ITO-coated glass as a blank reference. The measured refractive index was extracted from an experimentally fitted oscillator model based on a previous study [51].

Transmission spectra were taken using a commercial white light transmission microscopy setup (Witec Alpha series 300). The homogeneous metasurface sample was illuminated by a normal incident collimated white light with a linear polarization. The transmitted light was collected using a  $20 \times$  objective with NA = 0.4 and direct to a grating-based spectrometer.

Diffraction patterns of the transmitted light were collected on a home-built optical setup as shown in Figure S4. A white light source as well as a 785 nm laser was coupled into the optical path as light sources. The custom-built electrochemical cell with the metasurface sample was placed normally to the incident laser, where a  $20 \times$  objective with NA = 0.4 was used to collect the optical response of the fabricated nanostructures. A high-speed camera with a monochromatic CMOS device (Basler acA1440-220um) was employed to locate the sample and monitor the diffraction patterns.

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metasurface and carried out the nanofabrication. L.de.S.M. and W.L. constructed the optical measurement setup. W.L. designed the customized electrochemical cell and conducted the electrical switching experiments. A.T. and H.R. contributed to the metasurface design and made helpful comment on the manuscript. S.A.M. offered ideal facility support for carrying out the experiments. All authors discussed the results, analyzed the data and commented on the manuscript.

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**Data availability:** All the relevant data that support the findings of this study are available from the corresponding author upon reasonable request.

### References

- [1] T. Gu, H. J. Kim, C. Rivero-Baleine, and J. Hu, "Reconfigurable metasurfaces towards commercial success," *Nat. Photonics*, vol. 17, no. 1, pp. 48–58, 2023.
- [2] D. N. Neshev and A. E. Miroshnichenko, "Enabling smart vision with metasurfaces," *Nat. Photonics*, vol. 17, no. 1, pp. 26–35, 2023.
- [3] A. M. Shaltout, V. M. Shalaev, and M. L. Brongersma, "Spatiotemporal light control with active metasurfaces," *Science*, vol. 364, no. 6441, p. eaat3100, 2019.
- [4] J. Yu, et al., "Electrically tunable nonlinear polaritonic metasurface," Nat. Photonics, vol. 16, no. 1, pp. 72–78, 2022.
- [5] P. C. Wu, *et al.*, "Dynamic beam steering with all-dielectric electro-optic III-V multiple-quantum-well metasurfaces," *Nat. Commun.*, vol. 10, p. 3654, 2019.
- [6] J. Karst, *et al.*, "Electrically switchable metallic polymer nanoantennas," *Science*, vol. 374, no. 6567, pp. 612–616, 2021.
- [7] D. de Jong, *et al.*, "Electrically switchable metallic polymer metasurface device with gel polymer electrolyte," *Nanophotonics*, vol. 12, no. 8, pp. 1397–1404, 2023.
- [8] S. Abdollahramezani, *et al.*, "Electrically driven reprogrammable phase-change metasurface reaching 80% efficiency," *Nat. Commun.*, vol. 13, p. 1696, 2022.
- [9] S.-Q. Li, X. Xu, R. M. Veetil, V. Valuckas, R. Paniagua-Domínguez, and A. I. Kuznetsov, "Phase-only transmissive spatial light modulator based on tunable dielectric metasurface," *Science*, vol. 364, no. 6445, pp. 1087–1090, 2019.
- [10] J. Park, et al., "All-solid-state spatial light modulator with independent phase and amplitude control for three-dimensional LiDAR applications," *Nat. Nanotechnol.*, vol. 16, no. 1, pp. 69–76, 2021.

- [11] H. Kwon, T. Zheng, and A. Faraon, "Nano-electromechanical spatial light modulator enabled by asymmetric resonant dielectric metasurfaces," *Nat. Commun.*, vol. 13, p. 5811, 2022.
- G. K. Shirmanesh, R. Sokhoyan, P. C. Wu, and H. A. Atwater,
  "Electro-optically tunable multifunctional metasurfaces," *ACS Nano*, vol. 14, no. 6, pp. 6912–6920, 2020.
- [13] M. Y. Shalaginov, *et al.*, "Reconfigurable all-dielectric metalens with diffraction-limited performance," *Nat. Commun.*, vol. 12, p. 1225, 2021.
- [14] J. Karst, et al., "Electro-active metaobjective from metalenses-on-demand," Nat. Commun., vol. 13, p. 7183, 2022.
- [15] A. Afridi, J. Canet-Ferrer, L. Philippet, J. Osmond, P. Berto, and R. Quidant, "Electrically driven varifocal silicon metalens," ACS *Photonics*, vol. 5, no. 11, pp. 4497–4503, 2018.
- [16] R. Kaissner, *et al.*, "Electrochemically controlled metasurfaces with high-contrast switching at visible frequencies," *Sci. Adv.*, vol. 7, no. 19, p. eabd9450, 2021.
- [17] J. Li, P. Yu, S. Zhang, and N. Liu, "Electrically-controlled digital metasurface device for light projection displays," *Nat. Commun.*, vol. 11, p. 3574, 2020.
- [18] I. Kim, et al., "Pixelated bifunctional metasurface-driven dynamic vectorial holographic color prints for photonic security platform," *Nat. Commun.*, vol. 12, p. 3614, 2021.
- [19] J. Li, S. Kamin, G. Zheng, F. Neubrech, S. Zhang, and N. Liu, "Addressable metasurfaces for dynamic holography and optical information encryption," *Sci. Adv.*, vol. 4, no. 6, p. eaar6768, 2018.
- [20] X. Duan, S. Kamin, and N. Liu, "Dynamic plasmonic colour display," *Nat. Commun.*, vol. 8, p. 14606, 2017.
- [21] J. Li, Y. Chen, Y. Hu, H. Duan, and N. Liu, "Magnesium-based metasurfaces for dual-function switching between dynamic holography and dynamic color display," ACS Nano, vol. 14, no. 7, pp. 7892–7898, 2020.
- [22] H.-S. Ee and R. Agarwal, "Tunable metasurface and flat optical zoom lens on a stretchable substrate," *Nano Lett.*, vol. 16, no. 4, pp. 2818–2823, 2016.
- [23] A. L. Holsteen, A. F. Cihan, and M. L. Brongersma, "Temporal color mixing and dynamic beam shaping with silicon metasurfaces," *Science*, vol. 365, no. 6450, pp. 257–260, 2019.
- [24] Q. Li, et al., "Metasurface optofluidics for dynamic control of light fields," Nat. Nanotechnol., vol. 17, no. 10, pp. 1097–1103, 2022.
- [25] I.-C. Benea-Chelmus, M. L. Meretska, D. L. Elder, M. Tamagnone, L. R. Dalton, and F. Capasso, "Electro-optic spatial light modulator from an engineered organic layer," *Nat. Commun.*, vol. 12, p. 5928, 2021.
- [26] Y. Zhang, et al., "Electrically reconfigurable non-volatile metasurface using low-loss optical phase-change material," Nat. Nanotechnol., vol. 16, no. 6, pp. 661–666, 2021.
- [27] Y. Wang, et al., "Electrical tuning of phase-change antennas and metasurfaces," Nat. Nanotechnol., vol. 16, no. 6, pp. 667–672, 2021.
- [28] A. Komar, et al., "Dynamic beam switching by liquid crystal tunable dielectric metasurfaces," ACS Photonics, vol. 5, no. 5, pp. 1742–1748, 2018.
- [29] X. Liu, et al., "Thermally dependent dynamic meta-holography using a vanadium dioxide integrated metasurface," Adv. Opt. Mater., vol. 7, no. 12, p. 1900175, 2019.
- [30] X. Duan, *et al.*, "Reconfigurable multistate optical systems enabled by VO<sub>2</sub> phase transitions," *ACS Photonics*, vol. 7, no. 11, pp. 2958–2965, 2020.

- [31] P. Kepič, *et al.*, "Optically tunable Mie resonance VO<sub>2</sub> nanoantennas for metasurface in the visible," *ACS Photonics*, vol. 8, no. 4, pp. 1048–1057, 2021.
- [32] M. R. Shcherbakov, et al., "Ultrafast all-optical tuning of direct-gap semiconductor metasurfaces," Nat. Commun., vol. 8, p. 17, 2017.
- [33] A. Leitis, *et al.*, "All-dielectric programmable Huygens' metasurfaces," *Adv. Funct. Mater.*, vol. 30, no. 19, p. 1910259, 2020.
- [34] Q. Wang, *et al.*, "Optically reconfigurable metasurfaces and photonic devices based on phase change materials," *Nat. Photonics*, vol. 10, no. 1, pp. 60–65, 2016.
- [35] P. P. Iyer, M. Pendharkar, C. J. Palmstrøm, and J. A. Schuller, "III-V heterojunction platform for electrically reconfigurable dielectric metasurfaces," *ACS Photonics*, vol. 6, no. 6, pp. 1345–1350, 2019.
- [36] J. Lee, et al., "Ultrafast electrically tunable polaritonic metasurfaces," Adv. Opt. Mater., vol. 2, no. 11, pp. 1057–1063, 2014.
- [37] J. Eaves-Rathert, E. Kovalik, C. F. Ugwu, B. R. Rogers, C. L. Pint, and J. G. Valentine, "Dynamic color tuning with electrochemically actuated TiO<sub>2</sub> metasurfaces," *Nano Lett.*, vol. 22, no. 4, pp. 1626–1632, 2022.
- [38] Y. Li, J. van de Groep, A. A. Talin, and M. L. Brongersma, "Dynamic tuning of gap plasmon resonances using a solid-state electrochromic device," *Nano Lett.*, vol. 19, no. 11, pp. 7988–7995, 2019.
- [39] S. Chen and M. P. Jonsson, "Dynamic conducting polymer plasmonics and metasurfaces," ACS Photonics, vol. 10, no. 3, pp. 571–581, 2023.
- [40] W. Lu, N. Jiang, and J. Wang, "Active electrochemical plasmonic switching on polyaniline-coated gold nanocrystals," *Adv. Mater.*, vol. 29, no. 8, p. 1604862, 2017.
- W. Lu, T. H. Chow, S. N. Lai, B. Zheng, and J. Wang,
  "Electrochemical switching of plasmonic colors based on polyaniline-coated plasmonic nanocrystals," *ACS Appl. Mater. Interfaces*, vol. 12, no. 15, pp. 17733–17744, 2020.

- [42] Y. Lu, S. H. Lam, W. Lu, L. Shao, T. H. Chow, and J. Wang, "All-state switching of the Mie resonance of conductive polyaniline nanospheres," *Nano Lett.*, vol. 22, no. 3, pp. 1406–1414, 2022.
- [43] M. Decker, et al., "High-efficiency dielectric Huygens' surfaces," Adv. Opt. Mater., vol. 3, no. 6, pp. 813–820, 2015.
- [44] W. Lu, T. H. Chow, Y. Lu, and J. Wang, "Electrochemical coating of different conductive polymers on diverse plasmonic metal nanocrystals," *Nanoscale*, vol. 12, no. 42, pp. 21617–21623, 2020.
- [45] E. Pensa, J. Gargiulo, A. Lauri, S. Schlücker, E. Cortés, and S. A. Maier, "Spectral screening of the energy of hot holes over a particle plasmon resonance," *Nano Lett.*, vol. 19, no. 3, pp. 1867–1874, 2019.
- [46] A. A. Pud, "Stability and degradation of conducting polymers in electrochemical systems," *Synth. Met.*, vol. 66, no. 1, pp. 1–18, 1994.
- [47] J. C. Lacroix, K. K. Kanazawa, and A. Diaz, "Polyaniline: a very fast electrochromic material," *J. Electrochem. Soc.*, vol. 136, no. 5, pp. 1308–1313, 1989.
- [48] J. Heinze, B. A. Frontana-Uribe, and S. Ludwigs, "Electrochemistry of conducting polymers – persistent models and new concepts," *Chem. Rev.*, vol. 110, no. 8, pp. 4724–4771, 2010.
- [49] K. Xiong, O. Olsson, J. Svirelis, C. Palasingh, J. Baumberg, and A. Dahlin, "Video speed switching of plasmonic structural colors with high contrast and superior lifetime," *Adv. Mater.*, vol. 33, no. 41, p. 2103217, 2021.
- [50] E. D. Palik, Handbook of Optical Constants of Solids, vol. 3, USA, Academic Press, 1998.
- [51] H. A. Al-Attar, Q. H. Al-Alawina, and A. P. Monkman, "Spectroscopic ellipsometry of electrochemically prepared thin film polyaniline," *Thin Solid Films*, vol. 429, no. 1, pp. 286–294, 2003.

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